Dwarf Galaxy Formation with H<sub>2</sub>-regulated Star Formation

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# **The Missing Satellites Problem**

Reality

## **Dark Matter Simulation**



There is strong tension between the observed number of dwarf satellite galaxies and the predicted number of dark matter subhalos orbiting our Milky Way galaxy.



#### The Field Dwarf Galaxy Problem

Marchesini et al. (2009) [see also Fontanot et al. 2009, Cirasuolo et al. 2010]



Semi-analytic models of galaxy formation (including prescriptions for SN feedback!) over-predict the abundance of low mass galaxies and the stellar mass density at intermediate to high redshifts.





#### The Field Dwarf Galaxy Problem

## Hydrodynamical Galaxy Formation Simulations



# enzo

Astrophysical Adaptive Mesh Refinement http://code.google.com/p/enzo/

- Cosmological Adaptive Mesh Refinement
- Follows dark matter and hydrodynamics.
- Includes cooling, star formation, supernova feedback, etc.
- Community code
- I've been a contributing developer since 2005.

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- ≻ 12.5 Mpc box
- $\succ$  256<sup>3</sup> DM particles (3×10<sup>6</sup> M<sub> $\odot$ </sub>)
- > 256<sup>3</sup> root grid + 7 levels of AMR
- Self-consistent metal cooling
- H<sub>2</sub>-regulated star formation

#### "Standard" Star Formation Simulation



Krumholz & Tan (2007) model Constant SFR per free-fall time SF threshold:  $n_{\text{thresh}} = 50 \text{ cm}^{-3}$   $\dot{\rho}_{\text{SF}} = \epsilon_{\star} \frac{\rho_{\text{gas}}}{t_{\text{freefall}}} \propto \rho_{\text{gas}}^{3/2}$   $t_{\text{freefall}} = \sqrt{\frac{3\pi}{32 \, G \, \rho}}$  $\epsilon_{\star} = 0.01$ 

Kuhlen, Krumholz, Madau, Smith, Wise (2011, arXiv:1105.2376)

#### "Standard" Star Formation Simulation



Kuhlen, Krumholz, Madau, Smith, Wise (2011, submitted)

#### "Standard" Star Formation Simulation

Number Density [cm<sup>-3</sup>]

Stellar Age [yr]



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Only weak supernova feedback:
> Injection of thermal energy (ϵ=10<sup>-5</sup>) in central grid cell.
> No winds!

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#### How to suppress SF in low mass halos

The most commonly invoked mechanism to suppress star formation in low mass dark matter halos is **Supernova/Stellar Wind Feedback** and **UV Photoheating**.

#### 1) UV Photoheating

- Typically only effective below few x 10  $^9$   $M_{\odot}$  halos.
- Difficult to explain complicated SF histories if Milky Way dwarfs

#### 2) Supernova/Stellar Wind Feedback

- Undoubtedly plays an important role in nature!
- Its effectiveness in numerical simulations is very implementation dependent.
- Even hydro simulations with SN feedback have trouble matching observed stellar mass functions.
- In SAMs it typically just means a removal of some/all gas from the SF reservoir below some halo mass, or a halo-mass-dependent SF efficiency.

Is it the whole story? Are we just putting the answer we want in by hand? In my opinion other mechanisms should be considered... For example: Molecular Hydrogen Regulated Star Formation. cf. Gnedin et al. (2009), Gnedin & Kravtsov (2010, 2011)

Bigiel et al. (2008): observational Kennicutt-Schmidt relation from spatially resolved (< 1 kpc) **radio**, **IR**, **and UV** observations of 7 nearby spiral galaxies.



The star formation rate correlates better with molecular gas (H2) than with atomic gas (HI) surface density.

SFR correlates with H<sub>2</sub> even though it's not the primary coolant (CII, CO)!



Pelupessy et al. (2006), Robertson & Kravtsov (2008), Gnedin et al. (2009), Feldmann et al. (2010), Krumholz & Gnedin (2010) Make SFR proportional to  $\rho_{\rm H_2}$ :  $\dot{\rho}_{\rm SF} = \epsilon_\star \frac{\rho_{\rm H_2}}{t_{\rm freefall}} \propto f_{\rm H_2} \rho_{\rm gas}^{3/2}$ How to get  $f_{\rm H_2}$  during simulation runtime:

- 1) Full non-equilibrium chemistry with  $H_2$  formation on dust grains, coupled to radiation transfer with Lyman Werner shielding (e.g. Gnedin et al. 2009, Feldman et al. 2010).
- Use results from idealized 1-D RT calculations of H<sub>2</sub> formation-dissociation balance in giant atomic-molecular cloud complexes (KMT09: Krumholz, McKee, & Tumlinson (2008, 2009), McKee & Krumholz (2010)).

Radiative transfer: 
$$\hat{\mathbf{e}} \cdot \nabla I_{\nu} = -n \left(\frac{1}{2} f_{\mathrm{H}_{2}} \sigma_{\mathrm{H}_{2},\nu} + \sigma_{d,\nu}\right) I_{\nu}$$
  
H<sub>2</sub> formation-dissociation  
balance:  $f_{\mathrm{H}_{1},n^{2}} \mathcal{R} = \frac{f_{\mathrm{H}_{2}}}{2} n \int d\Omega \int_{\nu_{1}}^{\nu_{2}} d\nu \frac{I_{\nu}}{h_{\nu}} \sigma_{\mathrm{H}_{2},\nu} f_{\mathrm{diss},\nu}$   
Fully molecular  
 $\hat{\mathbf{f}}_{\mathrm{H}_{2}} = \frac{1 - \frac{3}{4} \frac{s}{1 + 0.25s}}{0.6 \tau_{c}}$   
LW-shielding opacity  
 $\tau_{c} = \Sigma_{\mathrm{HI}} / \mu_{H} Z' \sigma'_{d}$   
 $\chi = 71 \left(\frac{\sigma_{d,-21}}{\mathcal{R}_{-16.5}}\right) \frac{G'_{0}}{n_{\mathrm{H},0}}$   
FUV intensity in units of the Milky Way's,  $7.5 \times 10^{4} \mathrm{cm}^{3}$   
(Draine 1978)  
formation on dust grains  $\approx 1$ 

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With the assumption of 2-phase equilibrium between a Cold Neutral Medium and a Warm Neutral Medium, the minimum CNM density is proportional to the LW flux



$$n_{\rm min} \approx 31G'_0 \frac{Z'_d/Z'_g}{1+3.1(G'_0Z'_d/\zeta'_t)^{0.365}} \,\,{\rm cm}^{-3}$$

and the KMT09 prescription for  $f_{H_2}$  becomes independent of the LW intensity.

$$\chi = 2.3 \left( \frac{\sigma_{d,-21}}{\mathcal{R}_{-16.5}} \right) \frac{1 + 3.1 \left( Z/Z_{\rm SN} \right)^{0.365}}{\phi_{\rm CNM}}$$

Krumholz & Gnedin (2010): direct comparison between self-consistent cosmological simulations (ART) and KMT09 model at z=3.

#### Simulations:

- > Cosmological zoom-in simulations of 3 disk galaxies ( $Z/Z_{\odot}$ =0.5, 0.01, 0.18).
- > Non-equilibrium chemical network with  $H_2$  formation on dust (local Z).
- > Star formation, metal enrichment, and "live" radiation transfer of ionizing radiation.
- > LW shielding with Sobolev-like approximation:  $S_{\rm D} = e^{-D_{\rm MW}\sigma_0(n_{\rm HI} + 2n_{\rm H_2})L_{\rm Sob}}$







Make SFR proportional to  $\rho_{\rm H2}$ No SF density threshold!  $\dot{\rho}_{\rm SF} = \epsilon_{\star} \, \frac{\rho_{\rm H_2}}{t_{\rm freefall}} \propto f_{\rm H_2} \, \rho_{\rm gas}^{3/2}$  $\epsilon_{\star} = 0.01$  $f_{\rm H_2} \simeq 1 - \frac{3}{4} \frac{s}{1+0.25s}$  $s = \frac{\ln(1 + 0.6\chi + 0.01\chi^2)}{0.6\,\tau_c}$  $\chi = 0.77 \left( 1 + 3.1 \, Z^{\prime 0.365} \right)$  $\tau_c = \Sigma_{\rm HI} / \mu_H \, Z' \, \sigma'_d$ 

 $10^{-3} Z_{\circ}$  metallicity floor at z=10.

Further metal enrichment from SN injection: 0.25 M<sub>\*</sub>, yield=0.02.

	Make SFR proportional to $\rho_{\rm H_2}$						
2 2 2	)	z=4	SMGs		No S	SF density threshold! $\rho_{ m SF} = \epsilon_\star  rac{ ho_{ m H_2}}{t_{ m freefall}} \propto f_{ m H_2}   ho_{ m gas}^{3/2}$	
Name	$z_{\rm final}$	$ ho_{ m gas,SF}$	$n_{\rm thresh}$	$J_{\rm LW}/J_{\rm MW}$	[Z <sub>floor</sub> ]	Comment	
KT07 KT07_low KT07_high KMT09 KMT09_FLW1 KMT09_FLW10 KMT09_FLW100 KMT09_FLW1000 KMT09_FLW1000 KMT09_ZF4.0 KMT09_ZF2.5 KMT09_ZF2.0 KMT09_ZF210	$\begin{array}{c} 4.0 \\ 6.0 \\ 6.0 \\ 4.0 \\ 5.0 \\ 5.0 \\ 5.0 \\ 5.0 \\ 5.0 \\ 6.0 \\ 6.0 \\ 6.0 \\ 6.0 \\ 6.0 \end{array}$	$tot tot H_2 H_2 H_2 H_2 H_2 H_2 H_2 H_2 H_2 H_2$	$50  \mathrm{cm}^{-3}$ $5  \mathrm{cm}^{-3}$ $500  \mathrm{cm}^{-3}$ — — — — — — — — — — — — —	  1 10 100 1000  		$\begin{array}{c} \mbox{Krumholz \& Tan (2007) SF law} \\ \mbox{lower SF threshold} \\ \mbox{higher SF threshold} \\ \mbox{Krumholz et al. (2009): 2-phase equilibrium} \\ \mbox{one additional refinement level (maxlevel=8)} \\ \mbox{KMT09 with uniform LW} \\ \mbox{background of} \\ \mbox{increasing} \\ \mbox{intensity} \\ \mbox{lower Z}_{\rm floor} \\ \mbox{higher Z}_{\rm floor} \\ \mbox{even higher Z}_{\rm floor} \\ \mbox{at } z = 10 \end{array}$	
-4 10 <sup>-3</sup> Z <sub>o</sub> metallicity floor at z=10.							
$\begin{array}{ccccc} -1 & 0 & 1 & 2 & 3 & 4 \\ & & \log_{10} \Sigma_{\text{Gas}} \left[ M_{\odot} \text{ pc}^{-2} \right] \end{array}$					Furth injec	Further metal enrichment from SN injection: 0.25 M <sub>*</sub> , yield=0.02.	

## Comparisons with observational SF scaling laws

See also: Gnedin, Tassis, & Kravtsov (2009), Gnedin & Kravtsov (2010, 2011), Feldmann & Gnedin (2010)



The H<sub>2</sub>-regulated model reproduce the turnover in  $\Sigma_{\rm SFR}$  without an artificial density threshold.

The  $H_2$ -KS relation lies between the Genzel et al. (2010) z=0 - 3.5 relations for "normal" and "luminous mergers".

### **Metallicity Dependence**



see also Bolatto et al. (2011, arXiv:1107.1717)



Our model is able to capture the metallicity-dependence of the rollover in the KS relation.

 $H_2$  fractions as a function of total  $\Sigma_{gas}$  compare favorably with recent direct measurements in the SMC (Bolatto et al. 2011, arXiv:1107.1717).









Lower mass halos have lower star formation efficiency ( $f_* = M_*/M_{tot}$ ) owing to their **lower metallicity**.

Lower Z  $\Rightarrow$  Less Lyman-Werner shielding  $\Rightarrow$  Smaller  $f_{H_2} \Rightarrow$  Reduced star formation

#### Halo Mass Dependence of f<sub>\*</sub>

Low Mass Halos (M <  $10^{10}$  M<sub> $\odot$ </sub>)

High Mass Halos (M >  $10^{10} M_{\odot}$ )



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Without the 2-phase equilibrium assumption the  $f_*$ -suppression mass scale depends on the strength of the LW background.

[It also becomes dependent on a subgrid clumping factor, set to 30 here (Krumholz & Gnedin 2010).]







Observational luminosity functions from Bouwens et al. 2007, 2010.

Dust corrections very important! [Bouwens et al. 2010: 1.55, 0.625, 0.375, 0, 0 mags at z = 4, 5, 6, 7, 8.]

We calculate  $L_{UV}$  from SFR:  $L_{UV} = 8.0 \times 10^{27} \, (SFR/M_{\odot} \, yr^{-1}) \, erg \, s^{-1} \, Hz^{-1}$ 

Standard SF overpredicts LF. [except at z=4?]

H<sub>2</sub>-regulated SF improves agreement around sensitivity limit (MUV=-18).

H<sub>2</sub> suppression in this realization may be too strong for fainter systems.



Compares favorably with current (uncertain!) determinations utilizing ultra-deep rest-frame UV HST ACS/WFC3 observations coupled with stellar masses estimated from Spitzer rest-frame optical measurements. [Bouwens et al. 2009, 2010, Gonzalez et al. 2010, Labbé et al. 2009, 2010, Stark et al. 2009]

## Conclusions

- The are two dwarf galaxy problems in our understanding of the galaxy formation process:
  - 1) The Missing Satellites Problem
  - 2) The Field Dwarf Galaxy Problem
- Both are typically explained by invoking "supernova feedback", but other explanations should be considered. One example is H<sub>2</sub>-regulated star formation.
- Cosmological AMR hydrodynamical galaxy formation simulations with Enzo show that regulating SF by the H<sub>2</sub> abundance:
  - Reproduces the cutoff in  $\Sigma_{SFR}$  in the Kennicutt-Schmidt relation at ~10 M<sub>o</sub>/pc<sup>2</sup> without the need for a SF density threshold.
  - Matches the observed  $H_2$ -KS relation as reported by Genzel et al. (2010) at z=0-3.5.
  - Suppresses star formation in M <  $10^{10}$  M<sub> $\odot$ </sub> halos, because these galaxies aren't able to self-enrich as well as more massive halos.
  - Improves the agreement with (uncertain) observational determinations of the cosmic stellar mass density and SFR density evolution at z>4.
  - Helps to alleviate the dwarf galaxy problems.